

overlap contribution
(magnetic hyperfine interaction)

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overlap in the charge-charge interaction

What did we have for the charge-charge interaction:

order	multipole moment / field	first order quasi moment / quasi field	second order quasi moment / quasi field	...
O(0)	$\left. \begin{matrix} M \times r^0 Y_{00} \\ V \times r(0) \end{matrix} \right\} \text{MI } a $	$\left. \begin{matrix} M^{(1)} \times \{r^2 Y_{00}\} \\ \tilde{V}^{(1)} \times \Delta r(0) \end{matrix} \right\} \text{MS}^{(1)} a $	$\left. \begin{matrix} M^{(2)} \times \{r^4 Y_{00}\} \\ \tilde{V}^{(2)} \times \Delta^2 r(0) \end{matrix} \right\} \text{MS}^{(2)}$...
O(2)	$\left. \begin{matrix} Q \times r^2 Y_{20} \\ V_{ij} \times \partial_{ij} r(0) \end{matrix} \right\} \text{QI } b $	$\left. \begin{matrix} Q^{(1)} \times \{r^4 Y_{20}\} \\ \tilde{V}_i^{(1)} \times \partial_{ij} \Delta r(0) \end{matrix} \right\} \text{QS}^{(1)} b $	$\left. \begin{matrix} Q^{(2)} \times \{r^6 Y_{20}\} \\ \tilde{V}_{ij}^{(2)} \times \partial_{ij} \Delta^2 r(0) \end{matrix} \right\} \text{QS}^{(2)}$...
O(4)	$\left. \begin{matrix} H \times r^4 Y_{40} \\ V_{ijkl} \times \partial_{ijkl} r(0) \end{matrix} \right\} \text{HDI } c $	$\left. \begin{matrix} H^{(1)} \times \{r^6 Y_{40}\} \\ \tilde{V}_{ijkl}^{(1)} \times \partial_{ijkl} \Delta r(0) \end{matrix} \right\} \text{HDS}^{(1)}$	$\left. \begin{matrix} H^{(2)} \times \{r^8 Y_{40}\} \\ \tilde{V}_{ijkl}^{(2)} \times \partial_{ijkl} \Delta^2 r(0) \end{matrix} \right\} \text{HDS}^{(2)}$...
...

0th order contribution for a point nucleus

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...

first order correction to 0th order for overlap (or 'extended nucleus')

$$-\frac{eZ}{6\epsilon_0} \rho_c(\mathbf{0}) \langle r_n^2 \rangle$$

- vanishes if the nucleus is a point...
- ...even if the charge density at r=0 is not zero
- correction is small w.r.t. regular monopole term.

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overlap in the current-current interaction

Translate this to the leading term of the current-current interaction

	dipole contribution	first order correction due to overlap
O(1)	$-\vec{\mu}_I \cdot \vec{B}(0)$	

- orbital
- spin

- Fermi contact contribution
 - $-\frac{2\mu_B\mu_0}{3} (|\psi_{s,1}(0)|^2 - |\psi_{s,-1}(0)|^2)$
 - does not vanish if nucleus becomes a point
- field related to Bohr-Weisskopf effect

overlap ⇔ extended nucleus !

'correction' is often dominant w.r.t. regular dipole term (see bcc-Fe)

Bohr-Weisskopf effect

Two different isotopes of the same element have a different nuclear moment. The spatial distribution of this moment over the nuclear volume need not to be homogeneous:

$$E_{mag} = - \int_{nuc} \mathbf{B}_{hf} \cdot d\boldsymbol{\mu}_I$$

- 1) Bring the two naked nuclei in an externally applied B and measure the magnetic interaction energy. B is constant over the nuclear volume and goes out of the integral:

$$\frac{E_1}{E_2} = \frac{\mu_{I1}}{\mu_{I2}}$$
- 2) Bring both nuclei in the same hyperfine field. This varies over the nuclear volume and interacts with the (differently) distributed magnetic moments. B does not get out the integral, which leads to:

$$\frac{E_1}{E_2} = \frac{\mu_{I1}}{\mu_{I2}} (1 + \Delta)$$

→ Bohr-Weisskopf effect

↪ hyperfine anomaly (up to 2%)

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