

perturbation theory

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recapitulation: perturbation theory

Assuming we know eigenvalues and eigenfunctions of a hamiltonian H_0 :

$$H_0 |n_0\rangle = E_n^0 |n_0\rangle$$

what are the eigenvalues and eigenfunctions of a hamiltonian H that has the form

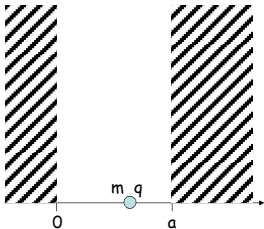
$$H = H_0 + \epsilon H_1 + \epsilon^2 H_2 + \dots$$

where ϵ is a small number ($\ll 1$) ?

(Note: H_2 can be zero).

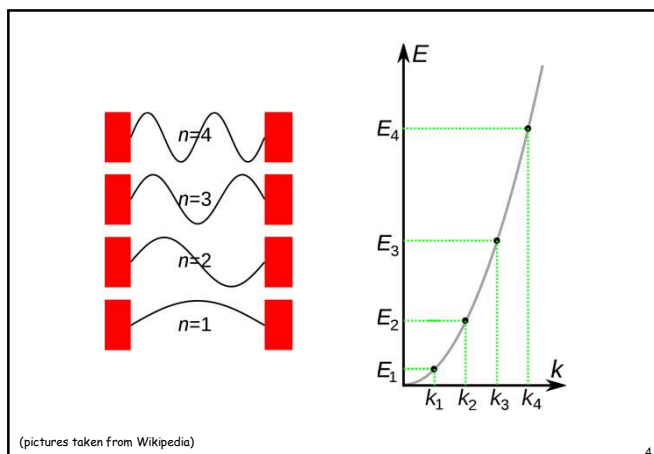
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Example 1: charged particle in an infinitely deep potential well



$$\begin{aligned} H_0 &= -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \\ E_{0,N} &= \frac{\hbar^2 \pi^2}{2ma^2} N^2 \quad (N = 1, 2, 3, \dots) \\ \Psi_{0,N}(x) &= \sqrt{\frac{2}{a}} \sin \frac{N\pi x}{a} \end{aligned}$$

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perturbation: constant electric field

electric field : $\vec{E}(x) = K\vec{e}_x$
 potential energy: $V(x) = -qKx$
 perturbing operator: $\hat{H}_1 = -qK\hat{x}$

$$\hat{H} = \hat{H}_0 + \hat{H}_1$$

$$= -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - qKx$$

form of hamiltonian $\propto \frac{d^2}{dx^2} + \underbrace{\frac{2qmK}{\hbar^2} x}_{\text{small if the electric field is not too strong}}$

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Solution up to first order in ϵ :

1) non-degenerate case

For eigenvalues that are non-degenerate (i.e. H_1 does not lift any degeneracy), the new eigenvalues and eigenfunctions are given by:

$$\Delta E_n^1 = \langle n_0 | H_1 | n_0 \rangle$$

$$\Rightarrow E_{\tilde{n}} = E_n^0 + \epsilon \Delta E_n^1$$

$$|\tilde{n}_1\rangle = \sum_{m \neq n} \frac{\langle m_0 | H_1 | n_0 \rangle}{E_n^0 - E_m^0} |m_0\rangle$$

$$\Rightarrow |\tilde{n}\rangle = |n_0\rangle + \epsilon |n_1\rangle$$

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Application to example 1:

$$\begin{aligned} \Delta E_N^1 &= \langle \Psi_{0,N} | \hat{H}_1 | \Psi_{0,N} \rangle \\ &= -\frac{2qK}{a} \int x \sin^2 \frac{N\pi x}{a} dx \\ &= -\frac{qKa}{2} \end{aligned}$$

Downward shift of all levels, independent of N.

Exercise: find the wave functions and probability density.

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Solution up to first order in ϵ :

2) l-fold degenerate case

For eigenvalues that are degenerate (i.e. H_1 does lift a degeneracy), the new eigenvalues and eigenfunctions are found by this procedure:

→ orthonormalize

Find an orthonormal basis $|n_0^l\rangle$ for the l-dimensional subspace

→ diagonalize

The l energy corrections are found as the l eigenvalues of this matrix:

$$\begin{bmatrix} \langle n_0^1 | \epsilon H_1 | n_0^1 \rangle & \langle n_0^1 | \epsilon H_1 | n_0^2 \rangle & \dots & \langle n_0^1 | \epsilon H_1 | n_0^l \rangle \\ \langle n_0^2 | \epsilon H_1 | n_0^1 \rangle & \langle n_0^2 | \epsilon H_1 | n_0^2 \rangle & \dots & \langle n_0^2 | \epsilon H_1 | n_0^l \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle n_0^l | \epsilon H_1 | n_0^1 \rangle & \langle n_0^l | \epsilon H_1 | n_0^2 \rangle & \dots & \langle n_0^l | \epsilon H_1 | n_0^l \rangle \end{bmatrix}$$

The new eigenstates are the eigenvectors of this matrix.

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Example 2: free electron under an applied magnetic field

without field, up and down spin are degenerate: $\Psi_\uparrow \Psi_\downarrow$

$$\begin{aligned} \hat{H}_0 &= \hat{\mathbb{I}} \\ \hat{H}_1 &= -\hat{\mu} \cdot \vec{B} \\ &= -\left(\frac{-2\mu_B \hat{S}}{\hbar}\right) \cdot \vec{B} \end{aligned}$$

$$\begin{bmatrix} \langle \Psi_\uparrow | \hat{S}_z | \Psi_\uparrow \rangle & \langle \Psi_\uparrow | \hat{S}_z | \Psi_\downarrow \rangle \\ \langle \Psi_\downarrow | \hat{S}_z | \Psi_\uparrow \rangle & \langle \Psi_\downarrow | \hat{S}_z | \Psi_\downarrow \rangle \end{bmatrix} = \begin{bmatrix} \frac{1}{2}\hbar & 0 \\ 0 & -\frac{1}{2}\hbar \end{bmatrix}$$

This is already diagonal.

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A derivation and more examples can be found at

[http://en.wikipedia.org/wiki/Perturbation_theory_\(quantum_mechanics\)](http://en.wikipedia.org/wiki/Perturbation_theory_(quantum_mechanics))

(section 2.1, 2.2 and 5)

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